

Appendix F

Statistical moments of total second order response

The total second order response can be represented in the following two forms:

$$X(t) = \int_{\tau} g_1(\tau)\zeta(t-\tau)d\tau + \int_{\tau_1} \int_{\tau_2} g_2(\tau_1, \tau_2)\zeta(t-\tau_1)\zeta(t-\tau_2)d\tau_1d\tau_2 \quad (\text{F.1})$$

or

$$= \sum_{i=1}^{\infty} (c_i + \lambda_i W_i) W_i \quad (\text{F.2})$$

From Eq.(F.2) the expected values up to third order are given as follows:

$$E[X] = \sum \lambda_i E[W_i^2] + \sum c_i E[W_i] \quad (\text{F.3})$$

$$E[X^2] = \sum c_i c_j E[W_i W_j] + \sum c_i \lambda_j E[W_i W_j^2] \quad (\text{F.4})$$

$$\begin{aligned} E[X^3] &= \sum c_i c_j c_k E[W_i W_j W_k] + \sum c_i \lambda_j c_k E[W_i W_j^2 W_k] \\ &+ \sum \lambda_i c_j c_k E[W_i^2 W_j W_k] + \sum \lambda_i \lambda_j c_k E[W_i^2 W_j^2 W_k] \\ &+ \sum c_i c_j \lambda_k E[W_i W_j W_k^2] + \sum c_i \lambda_j \lambda_k E[W_i W_j^2 W_k^2] \\ &+ \sum \lambda_i c_j \lambda_k E[W_i^2 W_j W_k^2] + \sum \lambda_i \lambda_j \lambda_k E[W_i^2 W_j^2 W_k^2] \end{aligned} \quad (\text{F.5})$$

Since $W_i (i = 1, \dots, \infty)$ are the standard Gaussian variables with mutual independence, the following relations⁷⁾ are satisfied:

$$E[W_i] = 0 \quad (\text{F.6})$$

$$E[W_i W_j] = \delta_{ij} \quad (\text{F.7})$$

$$E[W_i W_j W_k] = 0 \quad (\text{F.8})$$

$$E[W_i W_j W_k W_l] = \delta_{ij} \delta_{kl} + \delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk} \quad (\text{F.9})$$

$$E[W_i W_j W_k W_l W_m] = 0 \quad (\text{F.10})$$

$$\begin{aligned} E[W_i W_j W_k W_l W_m W_n] &= \delta_{ij} \delta_{kl} \delta_{mn} + \delta_{ij} \delta_{km} \delta_{ln} + \delta_{ij} \delta_{kn} \delta_{lm} \\ &\quad + \delta_{ik} \delta_{jl} \delta_{mn} + \delta_{ik} \delta_{jm} \delta_{ln} + \delta_{ik} \delta_{jn} \delta_{lm} \\ &\quad + \delta_{il} \delta_{jk} \delta_{mn} + \delta_{il} \delta_{jm} \delta_{kn} + \delta_{il} \delta_{jn} \delta_{km} \\ &\quad + \delta_{im} \delta_{jk} \delta_{ln} + \delta_{im} \delta_{jl} \delta_{kn} + \delta_{im} \delta_{jn} \delta_{kl} \\ &\quad + \delta_{in} \delta_{jk} \delta_{lm} + \delta_{in} \delta_{jl} \delta_{km} + \delta_{in} \delta_{jm} \delta_{kl} \end{aligned} \quad (\text{F.11})$$

for $i, j, k, l, m, n = 1, \dots, \infty$

where δ_{ij} is the Kronecker delta.

Using the above relations, the mean value \bar{X} or $E[X]$, the variance σ_X^2 or $Var[X]$, and the skewness ($\sqrt{\beta_1}$) are obtained as

$$\bar{X} = E[X] = \sum \lambda_i \quad (\text{F.12})$$

$$\sigma_X^2 = Var[X] = E[X^2] - \bar{X}^2 = \sum c_i^2 + 2 \sum \lambda_i^2 \quad (\text{F.13})$$

$$\sqrt{\beta_1} \sigma_X^3 = E[X^3] - 3E[X^2]E[X] + 2(E[X])^3 = 8 \sum \lambda_i^3 + 6 \sum c_i^2 \lambda_i \quad (\text{F.14})$$

While from Eq.(F.1) the expected values are written as:

$$E[X] = \int d\tau_1 \int d\tau_2 g_2(\tau_1, \tau_2) R_\zeta(\tau_2 - \tau_1) \quad (\text{F.15})$$

$$\begin{aligned} E[X^2] &= \int d\tau_1 \int d\tau_2 g_1(\tau_1) g_1(\tau_2) R_\zeta(\tau_2 - \tau_1) \\ &\quad + \int d\tau_1 \cdots \int d\tau_4 g_2(\tau_1, \tau_2) g_2(\tau_3, \tau_4) \\ &\quad \times [R_\zeta(\tau_2 - \tau_1) R_\zeta(\tau_4 - \tau_3) + R_\zeta(\tau_2 - \tau_3) R_\zeta(\tau_1 - \tau_4) \\ &\quad + R_\zeta(\tau_2 - \tau_4) R_\zeta(\tau_3 - \tau_1)] \end{aligned} \quad (\text{F.16})$$

$$\begin{aligned} E[X^3] &= \int d\tau_1 \cdots \int d\tau_4 g_2(\tau_1, \tau_2) g_1(\tau_3) g_1(\tau_4) \\ &\quad \times [6R_\zeta(\tau_3 - \tau_1) R_\zeta(\tau_1 - \tau_4) + 3R_\zeta(\tau_2 - \tau_1) R_\zeta(\tau_4 - \tau_3)] \\ &\quad + \left[\int d\tau_1 \int d\tau_2 g_2(\tau_1, \tau_2) R_\zeta(\tau_2 - \tau_1) \right]^3 \\ &\quad + \int d\tau_1 \cdots \int d\tau_6 g_2(\tau_1, \tau_2) g_2(\tau_3, \tau_4) g_2(\tau_5, \tau_6) \\ &\quad \times [6R_\zeta(\tau_2 - \tau_1) R_\zeta(\tau_6 - \tau_3) R_\zeta(\tau_5 - \tau_4) \\ &\quad + 8R_\zeta(\tau_3 - \tau_1) R_\zeta(\tau_5 - \tau_2) R_\zeta(\tau_6 - \tau_4)] \end{aligned} \quad (\text{F.17})$$

Transforming Eqs.(F.15), (F.16) and (F.17) to frequency domain we get:

$$\bar{X} = \int d\omega G_2(\omega, -\omega) S_\zeta(\omega) \quad (\text{F.18})$$

$$\begin{aligned} \sigma_X^2 &= \int d\omega |G_1(\omega)|^2 S_\zeta(\omega) \\ &+ \int d\omega_1 \int d\omega_2 |G_2(\omega_1, \omega_2)|^2 S_\zeta(\omega_1) S_\zeta(\omega_2) \end{aligned} \quad (\text{F.19})$$

$$\begin{aligned} \sqrt{\beta_1} \sigma_X^3 &= 6 \int d\omega_1 \int d\omega_2 G_1(-\omega_1) G_1(\omega_2) G_2(\omega_1, \omega_2) S_\zeta(\omega_1) S_\zeta(\omega_2) \\ &+ 8 \int d\omega_1 \int d\omega_2 \int d\omega_3 G_2(\omega_1, \omega_2) G_2^*(\omega_2, \omega_3) G_2(\omega_3, -\omega_1) \\ &\times S_\zeta(\omega_1) S_\zeta(\omega_2) S_\zeta(\omega_3) \end{aligned} \quad (\text{F.20})$$

where * denotes the complex conjugate.

Cross and Auto spectra

Take the cross correlation function between the nonlinear response process $X(t)$ and the Gaussian wave process $\zeta(t)$ as follows:

$$\begin{aligned} E[(X(t) - \bar{X})\zeta(t - \tau)] &= \int d\tau_1 g_1(\tau_1) E[\zeta(t - \tau_1)\zeta(t - \tau)] \\ &+ \int d\tau_1 \int d\tau_2 g_2(\tau_1, \tau_2) E[\zeta(t - \tau_1)\zeta(t - \tau_2)\zeta(t - \tau)] \\ &- \bar{X} E[\zeta(t - \tau)] \end{aligned} \quad (\text{F.21})$$

Since the wave process is defined to be zero-mean, the last two terms are zero. Thus:

$$E[(X(t) - \bar{X})\zeta(t - \tau)] = \int d\tau_1 g_1(\tau_1) E[\zeta(t - \tau_1)\zeta(t - \tau)] \quad (\text{F.22})$$

This means that the cross spectrum involves only the first term in the functional series (F.1), and thus that the linear transfer function G_1 is derivable by standard cross spectrum technique by Fourier transform. Denoting the cross spectrum as $S_{X\zeta}(\omega)$; then we get

$$S_{X\zeta}(\omega) = G_1(\omega) \cdot S_\zeta(\omega) \quad (\text{F.23})$$

Next, taking the auto correlation function of $X(t)$:

$$\begin{aligned} &E[(X(t) - \bar{X})(X(t + \tau) - \bar{X})] \\ &= \int d\tau_1 \int d\tau_2 g_1(\tau_1) g_1(\tau_2) E[\zeta(t - \tau_1)\zeta(t + \tau - \tau_2)] \end{aligned}$$

$$\begin{aligned}
& + \int d\tau_1 \cdots \int d\tau_4 g_2(\tau_1, \tau_2) g_2(\tau_3, \tau_4) \\
& \times E[\zeta(t - \tau_1) \zeta(t - \tau_2) \zeta(t + \tau - \tau_3) \zeta(t + \tau - \tau_4)] \\
& - \bar{X}^2
\end{aligned} \tag{F.24}$$

and using the factorization relation⁷⁾ for higher order moments of Gaussian processes as:

$$E[X_1 X_2 X_3 X_4] = E[X_1 X_2] E[X_3 X_4] + E[X_1 X_3] E[X_2 X_4] + E[X_1 X_4] E[X_2 X_3] \tag{F.25}$$

we obtain

$$\begin{aligned}
R_{XX}(\tau) & = \int d\tau_1 \int d\tau_2 g_1(\tau_1) g_1(\tau_2) R_\zeta(\tau + \tau_1 - \tau_2) \\
& + \int d\tau_1 \cdots \int d\tau_4 g_2(\tau_1, \tau_2) g_2(\tau_3, \tau_4) \\
& \times [R_\zeta(\tau + \tau_1 - \tau_3) R_\zeta(\tau + \tau_2 - \tau_4) \\
& + R_\zeta(\tau + \tau_1 - \tau_4) R_\zeta(\tau + \tau_2 - \tau_3)]
\end{aligned} \tag{F.26}$$

The auto power spectrum is the Fourier transform of R_{XX} and is computed from the Wiener-Khintchine relations as

$$S_X(\omega) = |G_1(\omega)|^2 S_\zeta(\omega) + 2 \int d\nu |G_2(\omega - \nu, \nu)|^2 S_\zeta(\omega - \nu) S_\zeta(\nu) \tag{F.27}$$

REFERENCES IN APPENDICES

- [1] Ogilvie, T. F. : Second-order hydrodynamic effects on ocean platforms, Proc. International Workshop on Ship and Platform Motions, 1983.
- [2] Dalzell, J.F.: Cross-Bi-Spectral Analysis : Application to Ship Resistance in Waves, J.S.R., vol.18, 1974.
- [3] Standing, R. G., Dacunha, N. M. C., Matten, R. B.: Mean Wave Drift Forces: Theory and Experiment, NMI R124, 1981.
- [4] Watson, G.N.: A Treatise on the Theory of Bessel Functions, Cambridge Univ. Press, 1966.
- [5] Kac, M. and Siegert, A.J.F.: On the Theory of Noise in Radio Receiver with Square-Law Detectors, Journ. Appl. Phys., vol.18, 1946.
- [6] Courant, R., and Hilbert, D.: Method of Mathematical Physics I, II, Interscience Pub., 1962.
- [7] Cramér, H.: Mathematical Method of Statistics, Princeton Univ. Press, 1946.

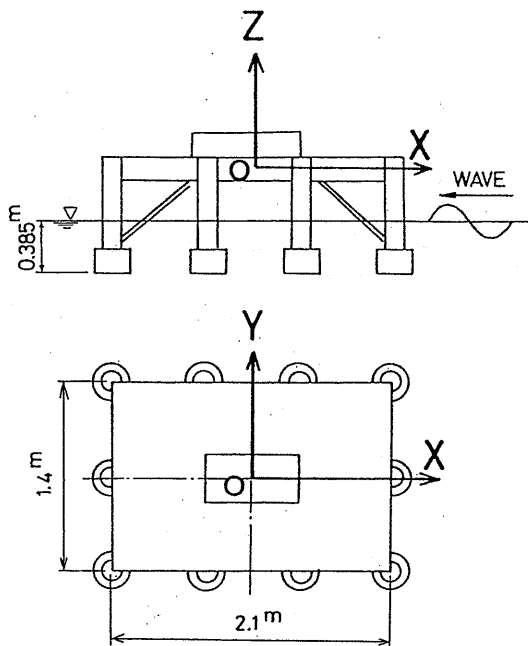


Figure 3.1 Configuration of a floating structure and the direction of incident waves

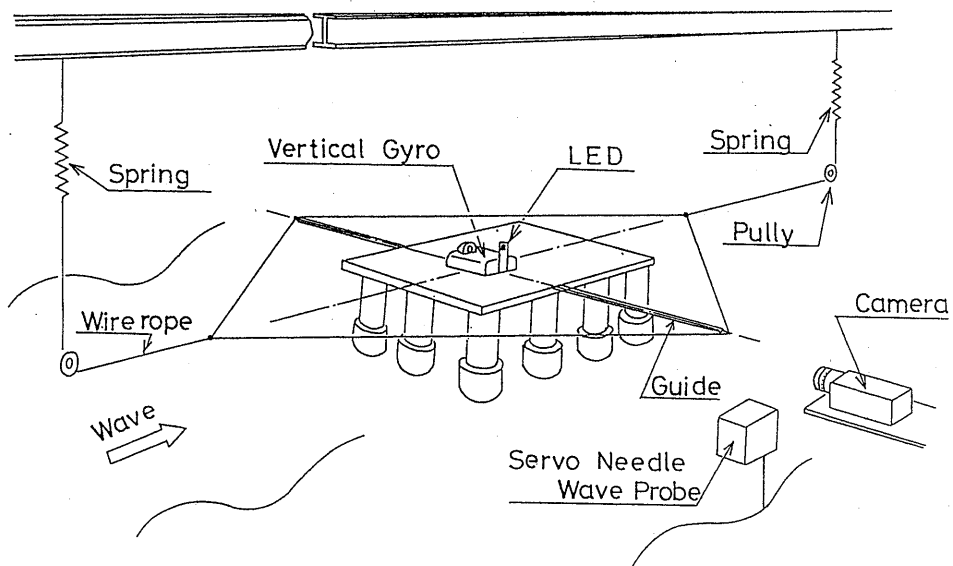


Figure 3.2 Set-up of model test

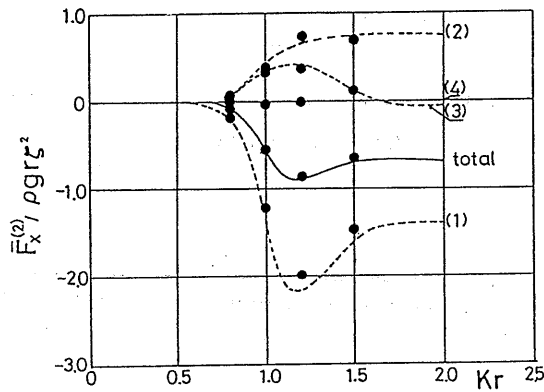
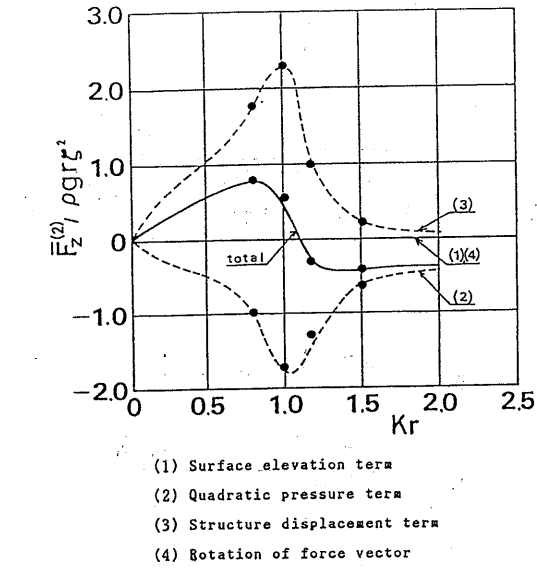


Figure 3.3 Components of the computed mean second order forces of a half sphere

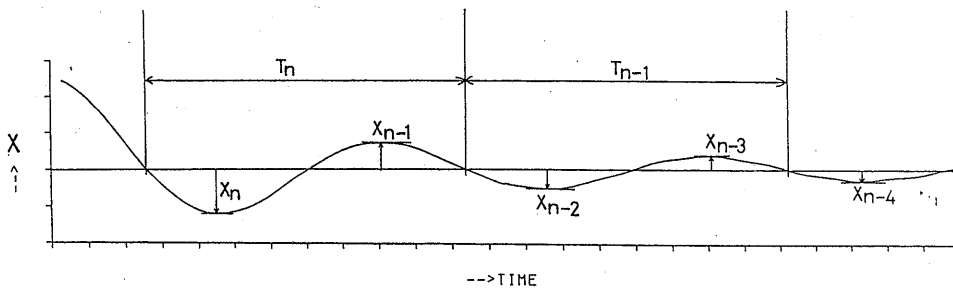


Figure 3.4 An example of the surge decaying motion

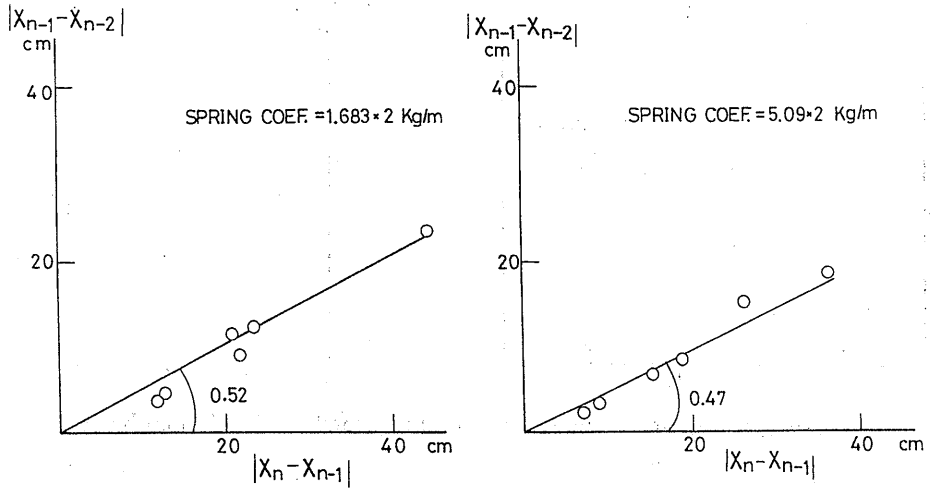


Figure 3.5 Extinction curves of surge motion

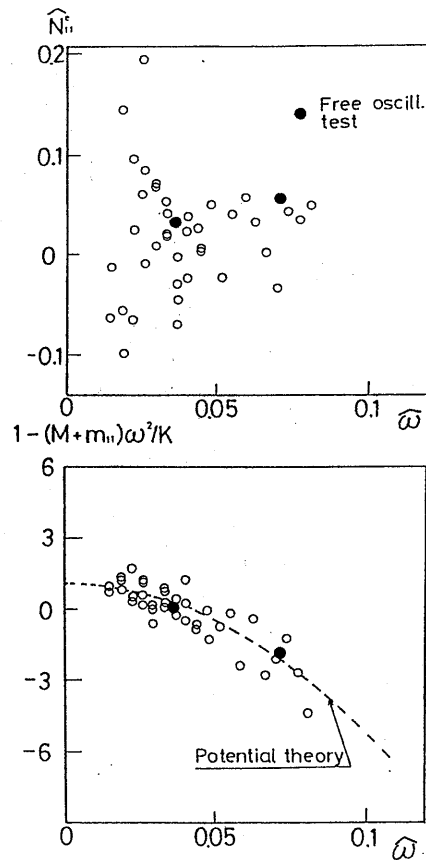


Figure 3.6 Added mass and damping coefficients (frequency base)

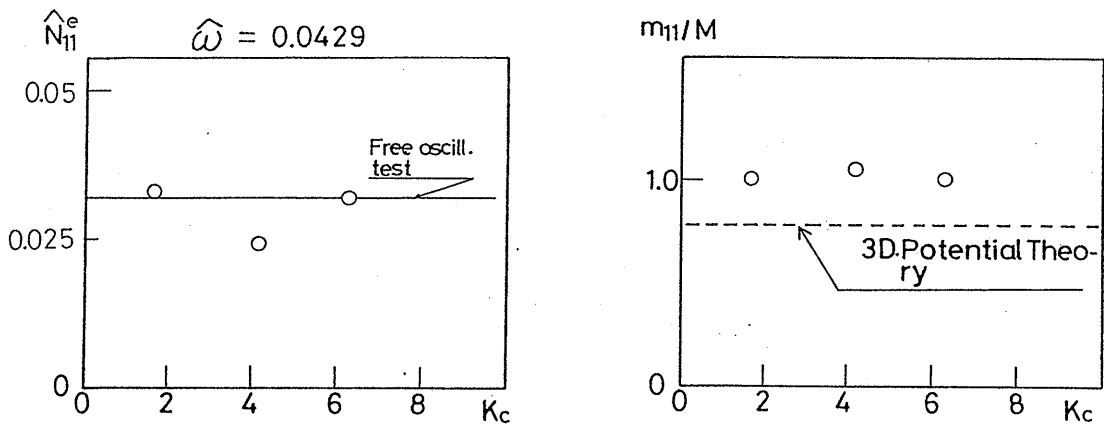


Figure 3.7 Added mass and damping coefficients (K_c number base)

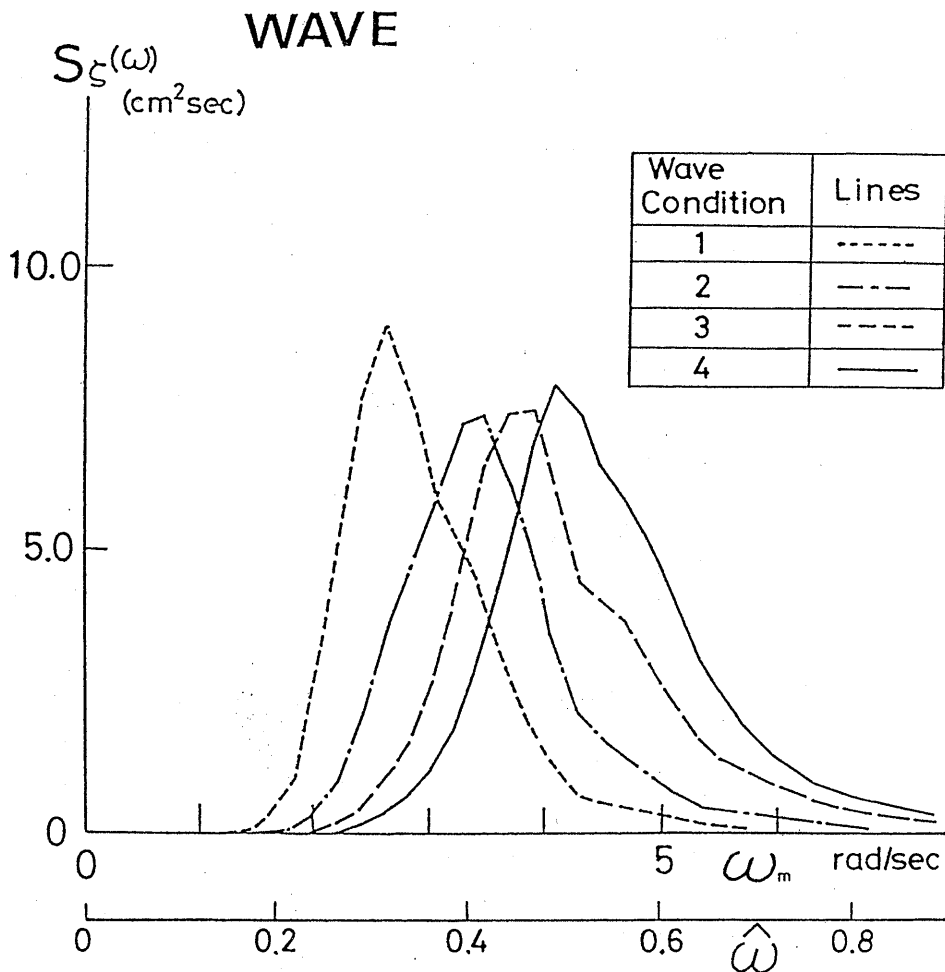


Figure 3.8 Wave spectra

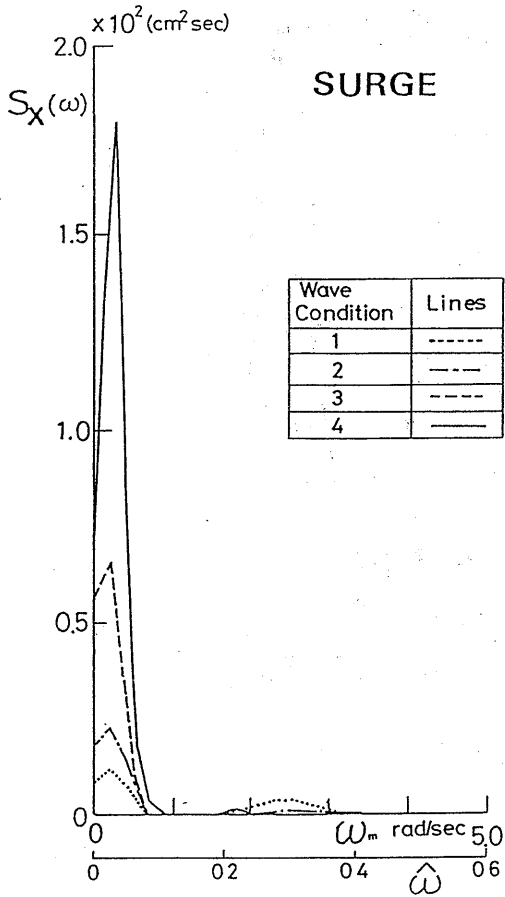


Figure 3.9 Surge response spectra

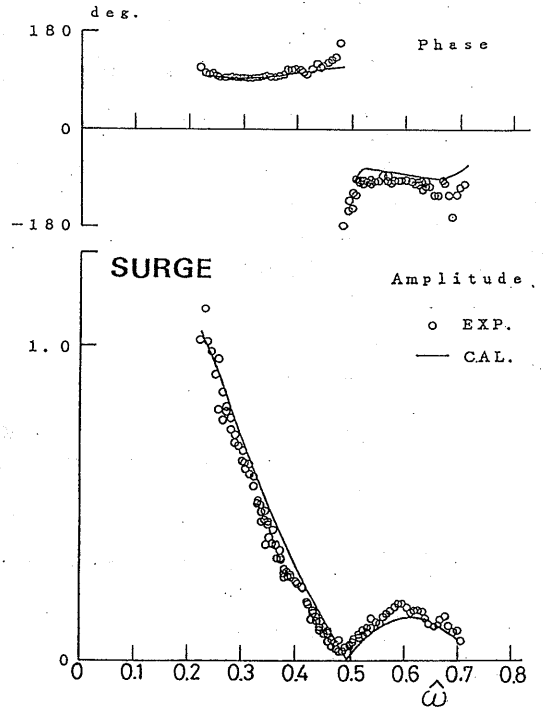


Figure 3.10 Linear transfer function $G_1(\omega)$ of surge motion

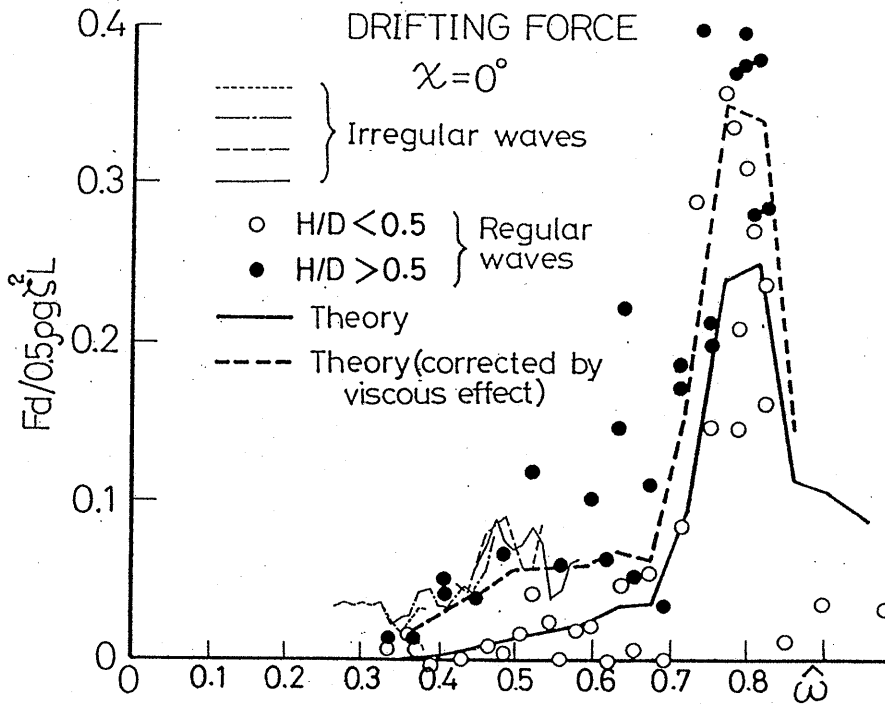


Figure 3.11 Longitudinal steady drift force in head waves

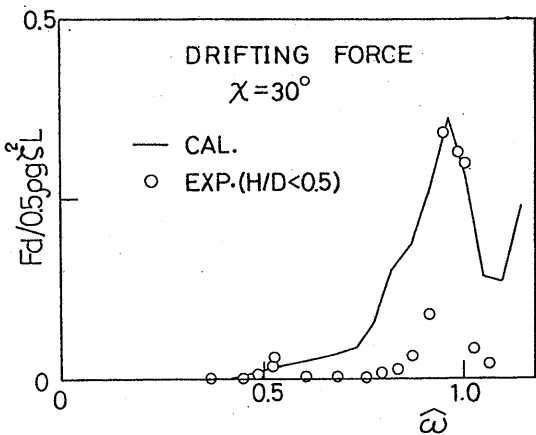


Figure 3.12 Longitudinal steady drift force in oblique waves ($\chi = 30^\circ$)

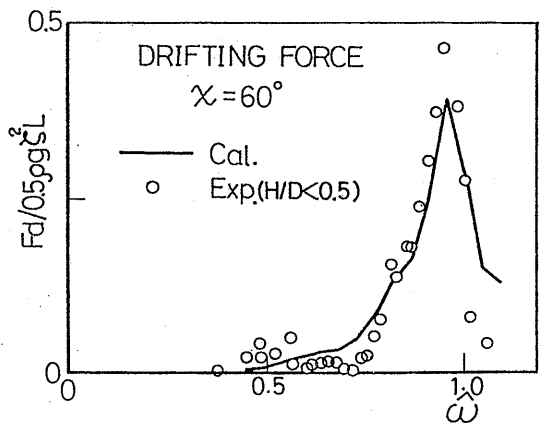


Figure 3.13 Longitudinal steady drift force in oblique waves ($\chi = 60^\circ$)

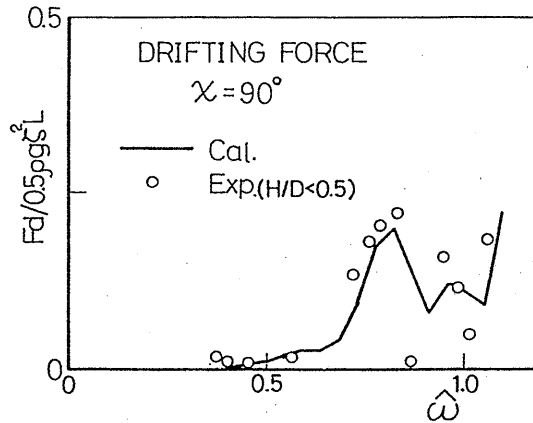


Figure 3.14 Longitudinal steady drift force in beam waves ($\chi = 90^\circ$)

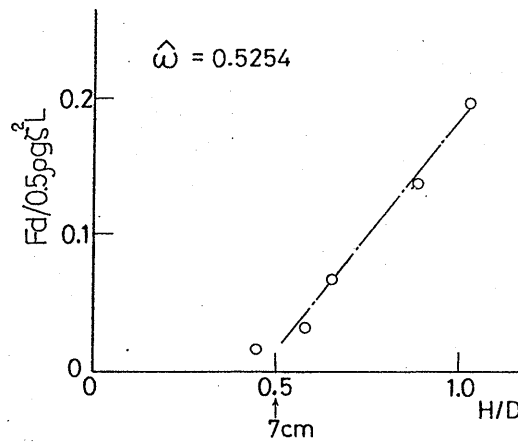


Figure 3.15 a) An example of wave height effect to steady drift force in head waves ($\hat{\omega} = 0.5254$)

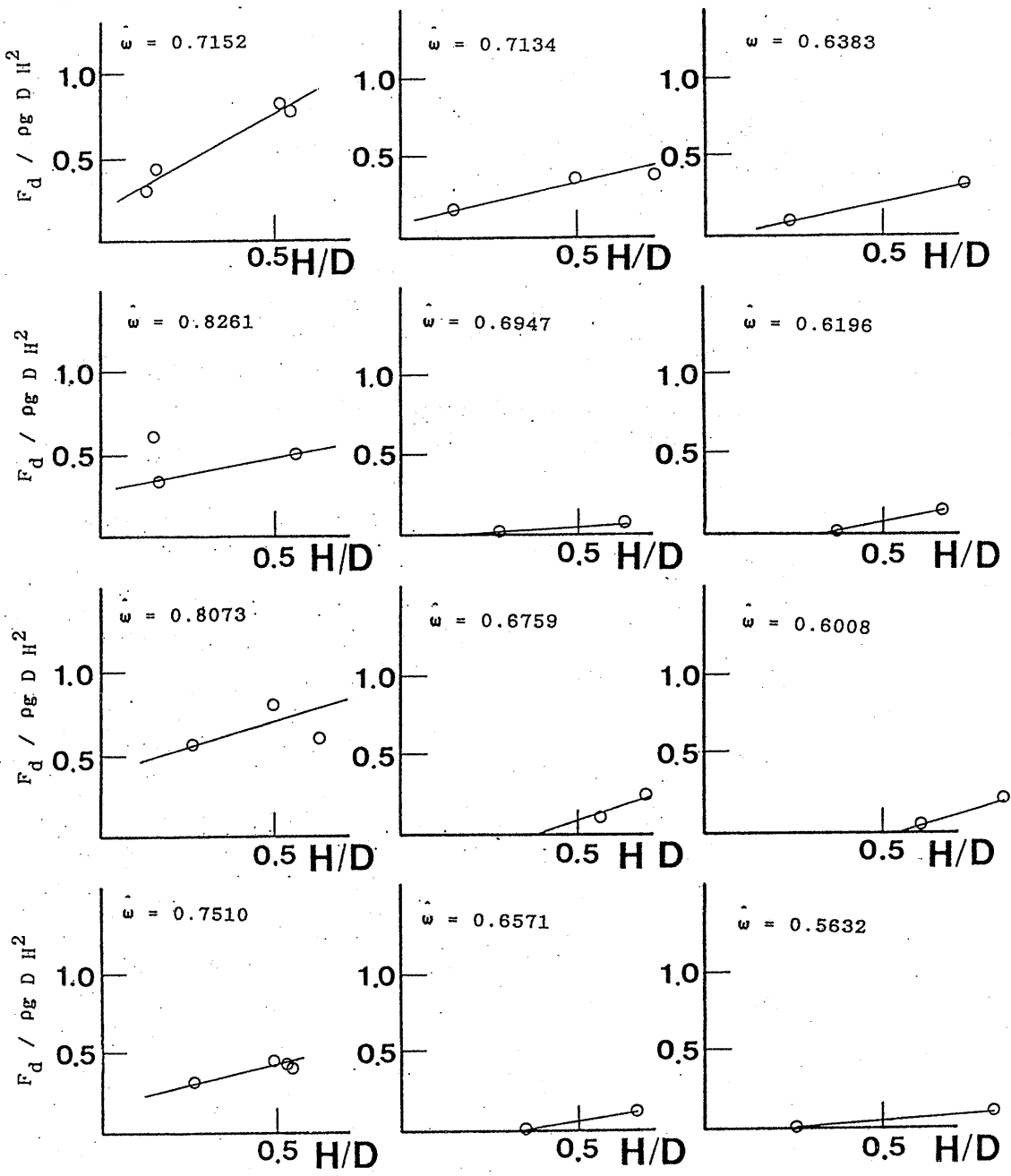


Figure 3.15 b) Wave height effects to steady drift force in head waves

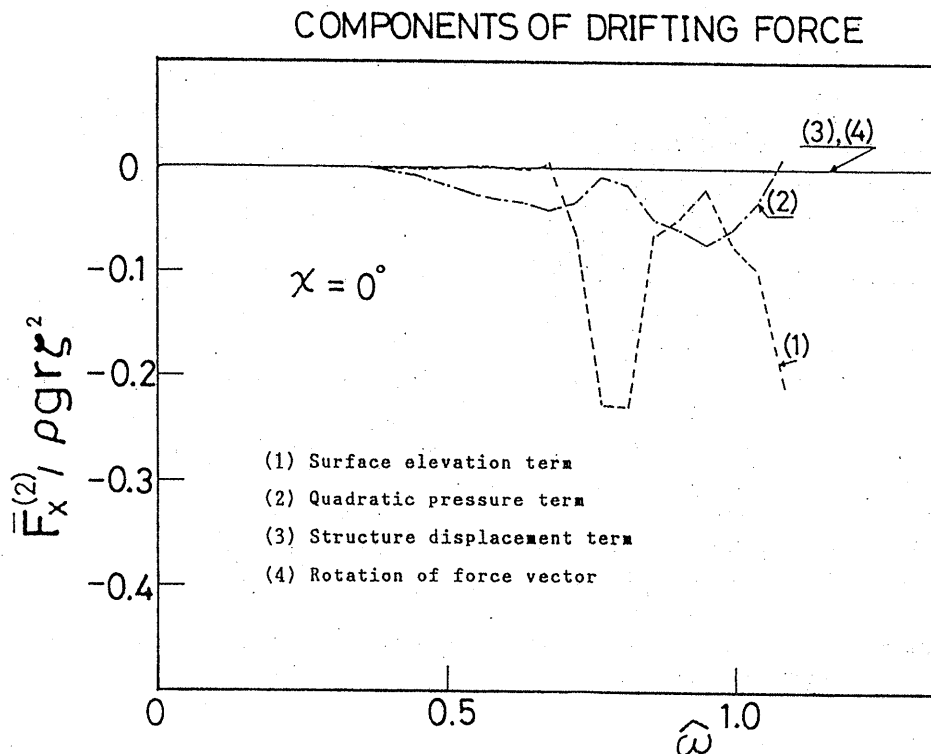


Figure 3.16 Components of steady drift force in head waves where r is half of the length L

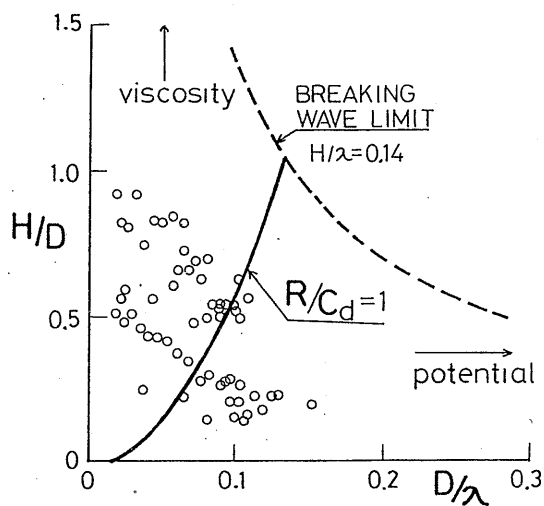


Figure 3.17 Comparison with viscous and potential components of drift forces to a vertical cylinder

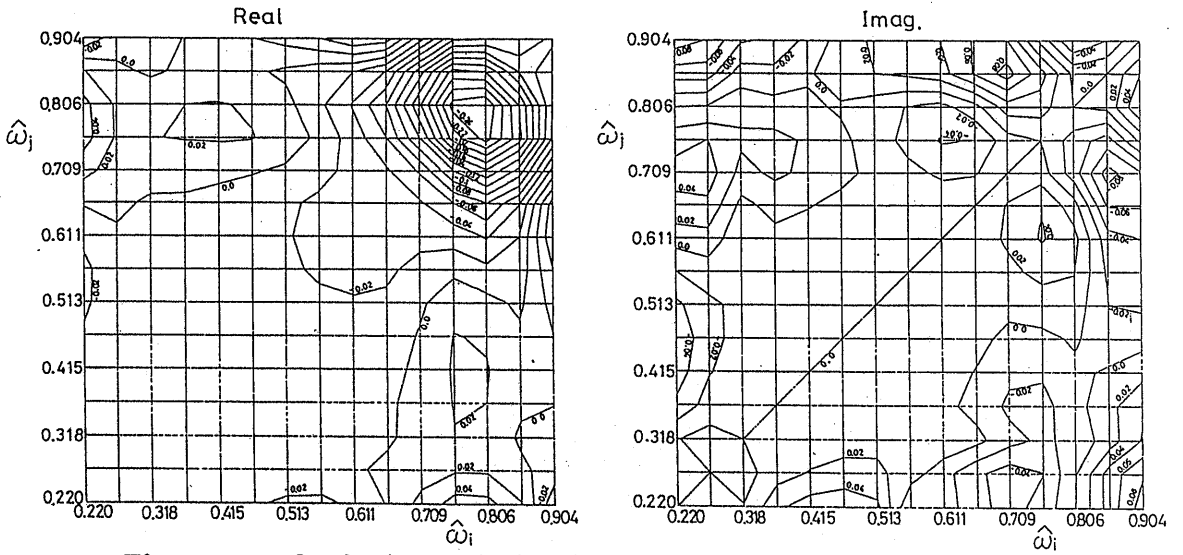


Figure 3.18 Quadratic transfer function of slowly varying drift force obtained from numerical calculation

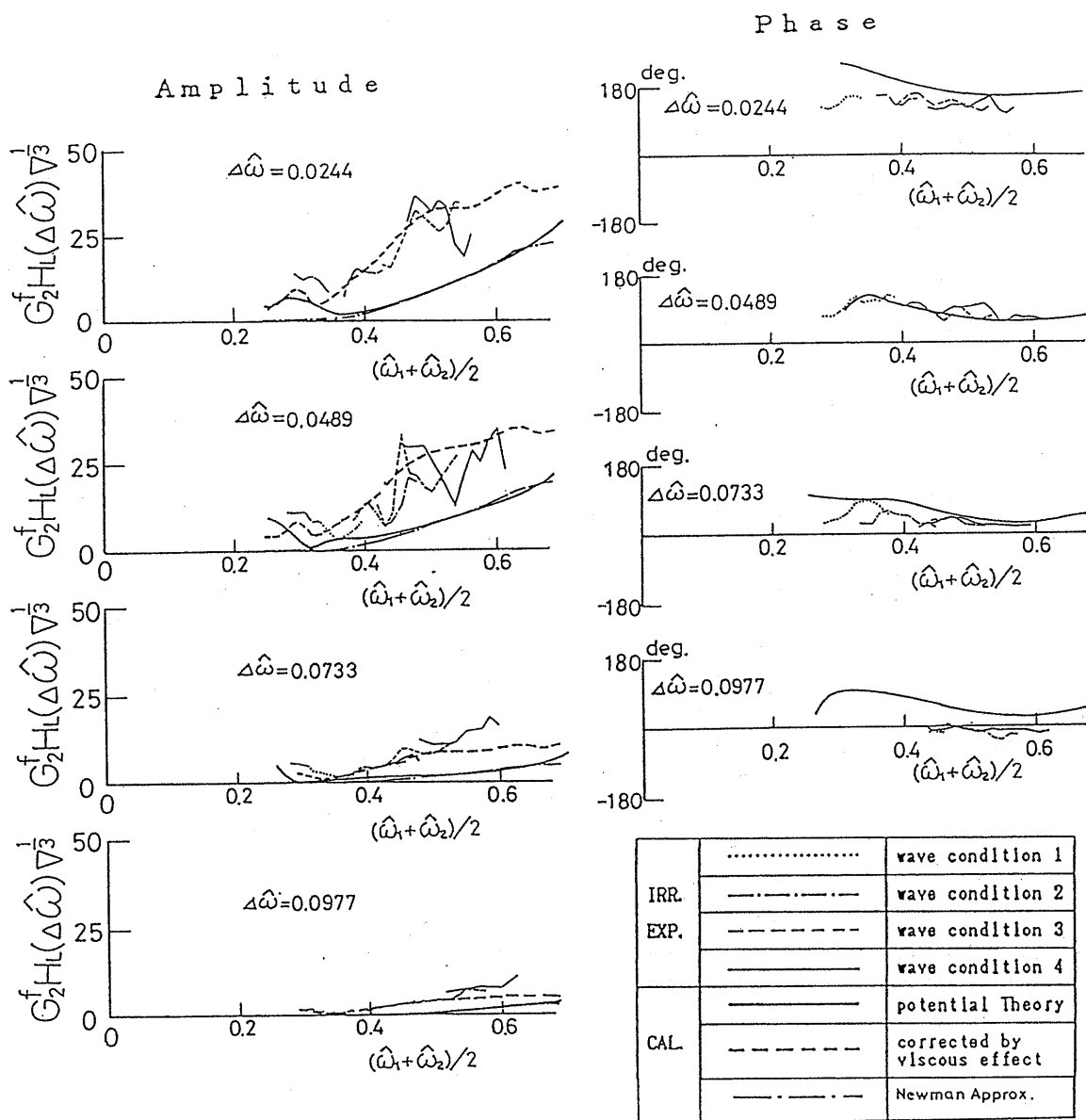


Figure 3.19 Comparison with experimental and numerical results on the quadratic transfer function of slowly varying drift force

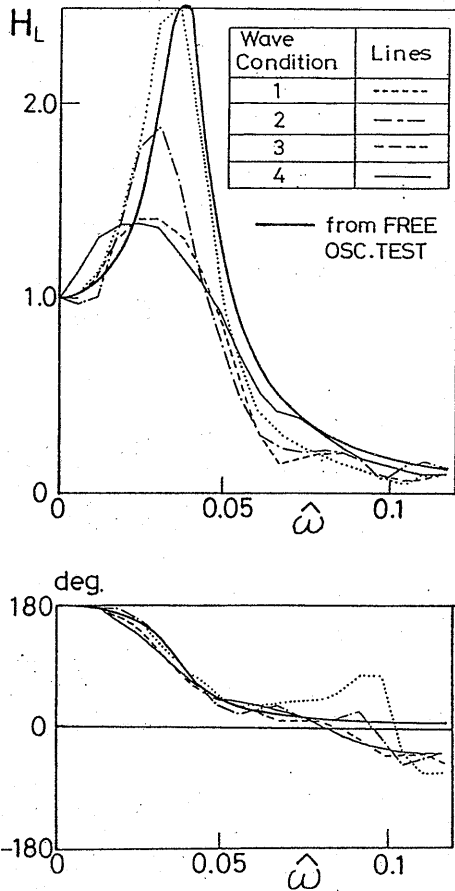


Figure 3.20 Frequency response function $H_L(\omega)$ of surge motion to external forces

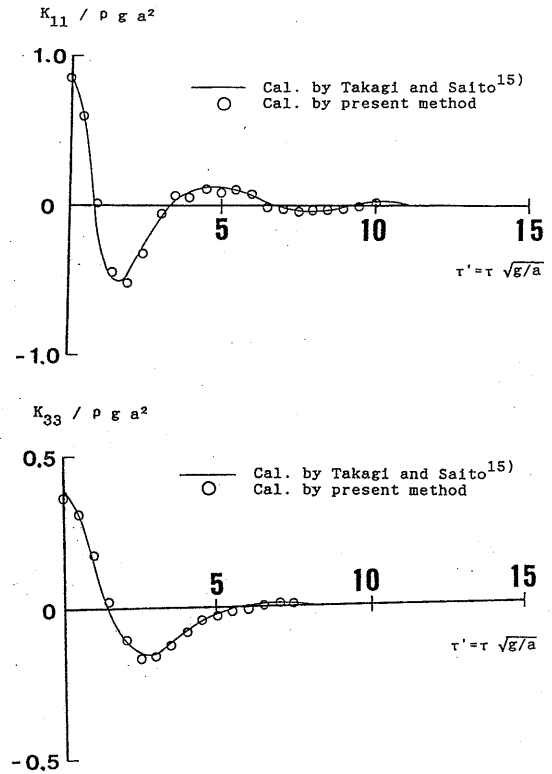


Figure 3.21 Memory effect functions of a half submerged sphere

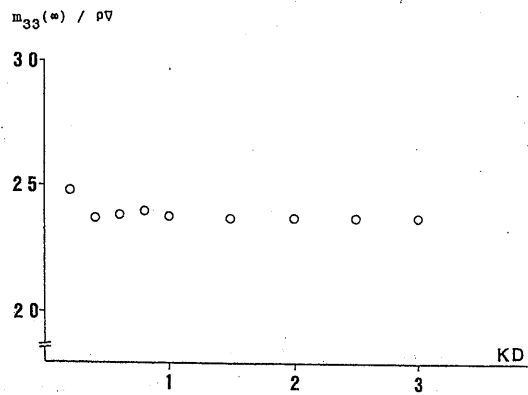


Figure 3.22 Heaving added mass at $\omega = \infty$ of a half submerged rectangular cylinder

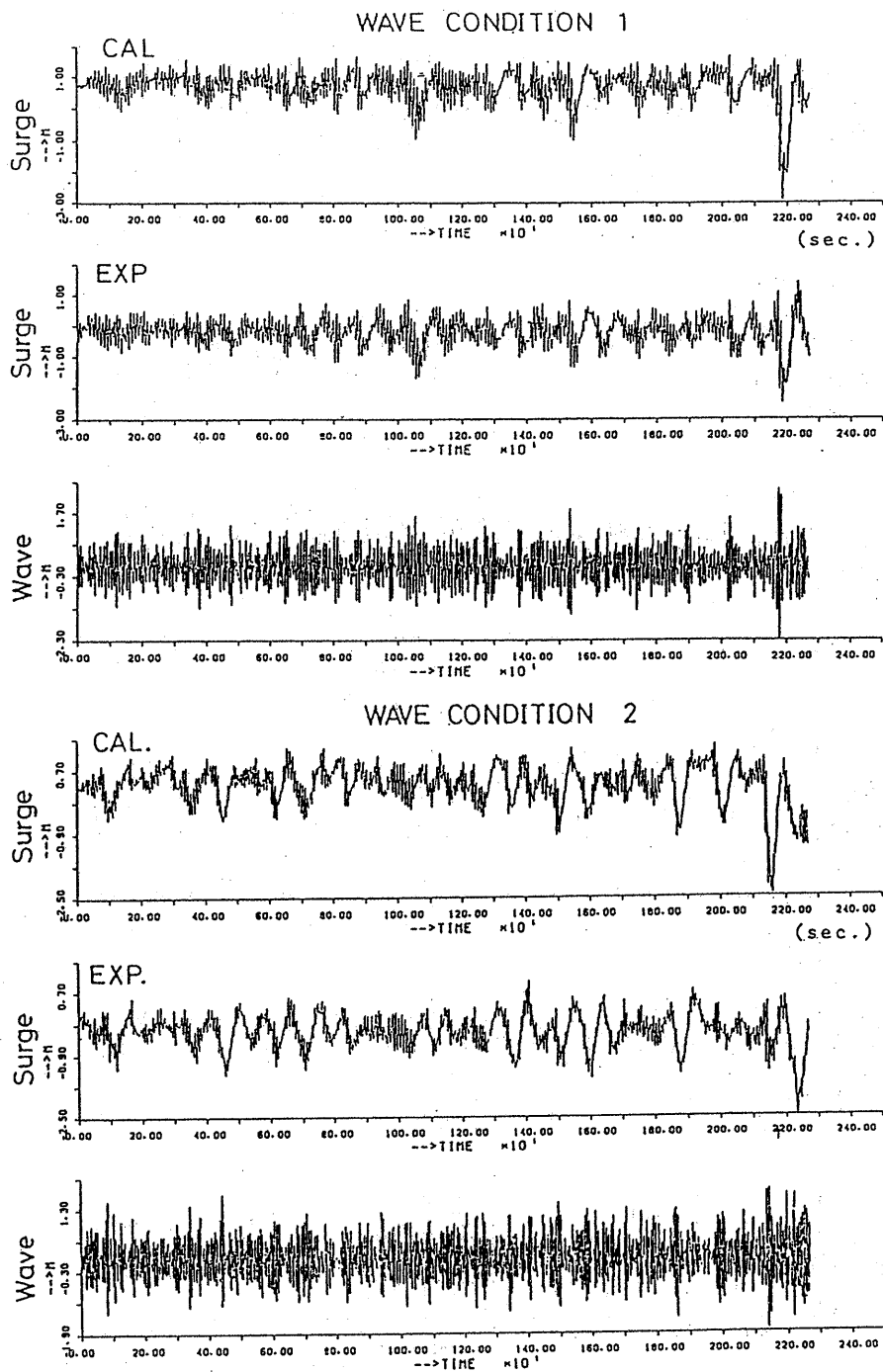


Figure 3.23 Comparisons between surge simulation results and measured ones on the floating body which is moored by linear springs (Wave conditions No1. and No2.)

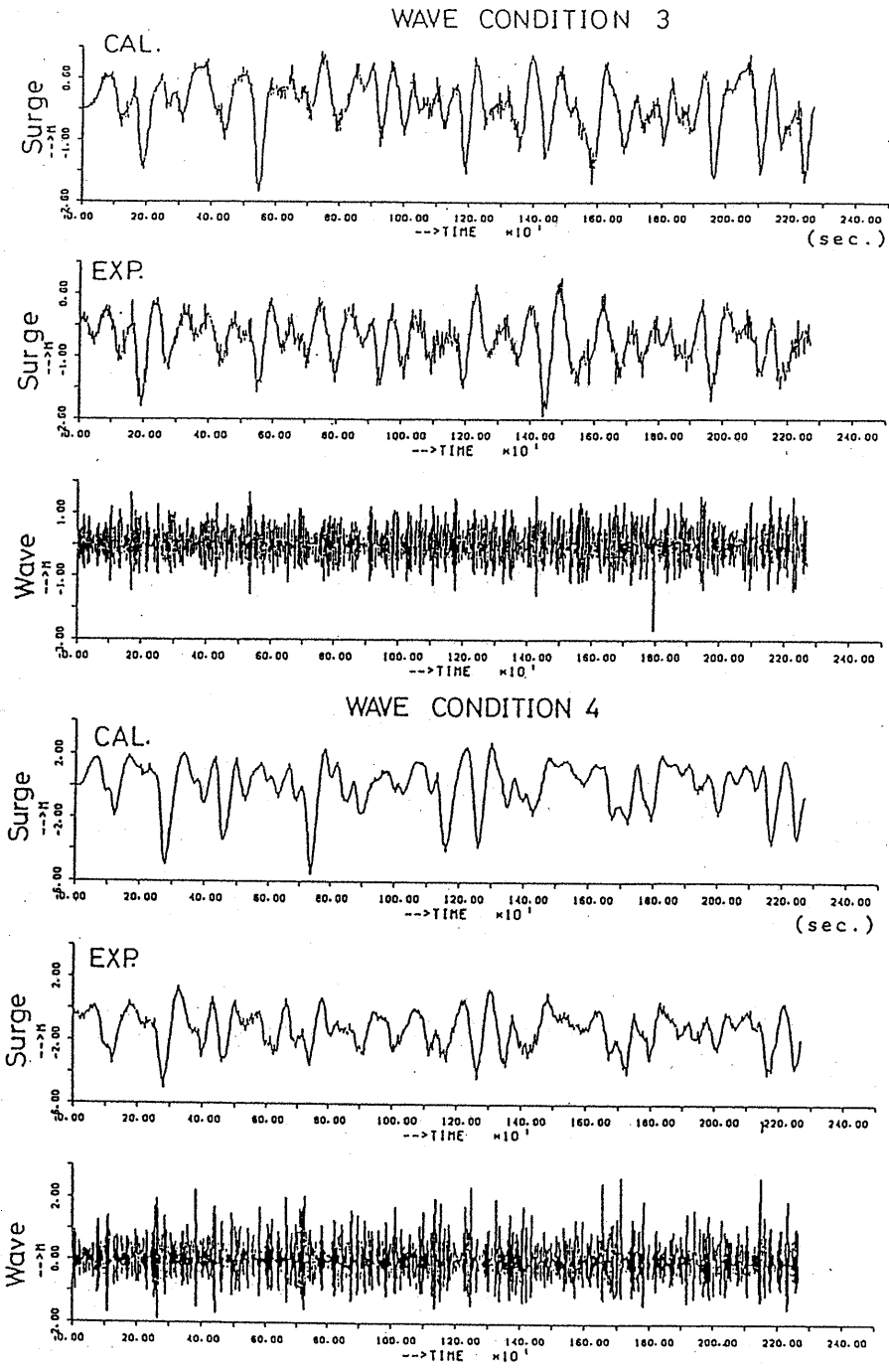


Figure 3.24 Comparisons between surge simulation results and measured ones on the floating body which is moored by linear springs (Wave conditions No.3. and No.4.)

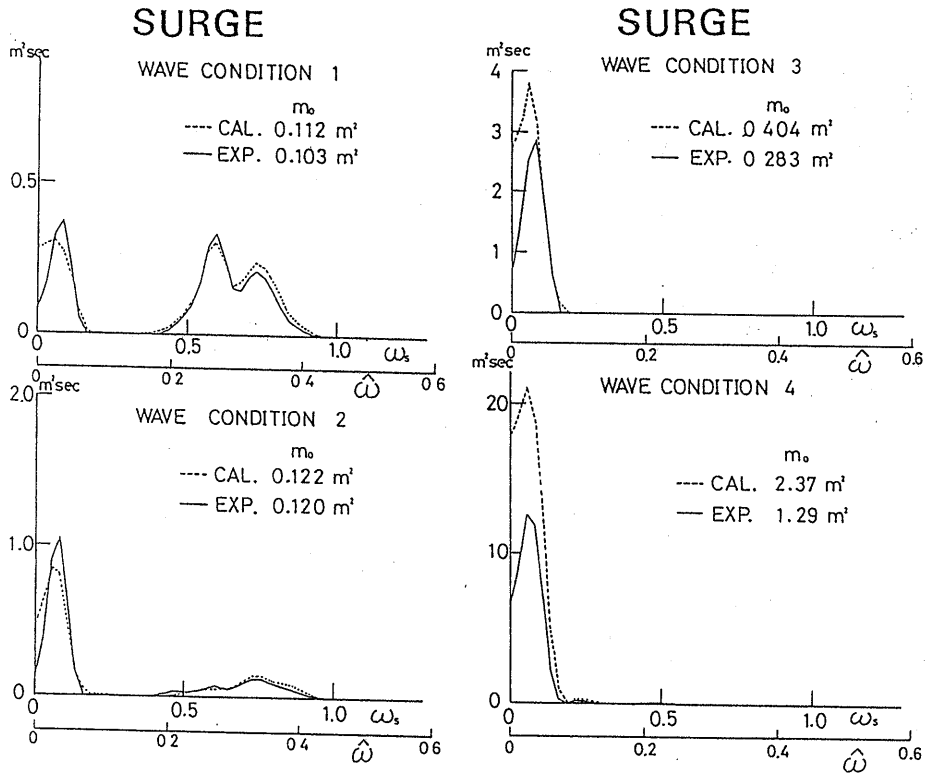


Figure 3.25 Comparisons with surge spectra of simulations and experiments

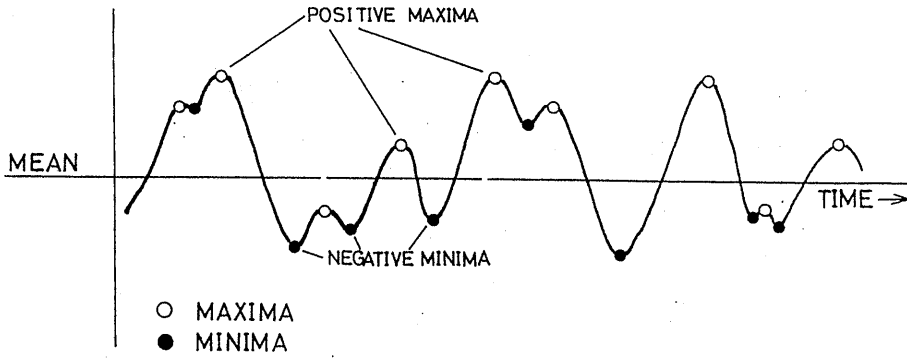
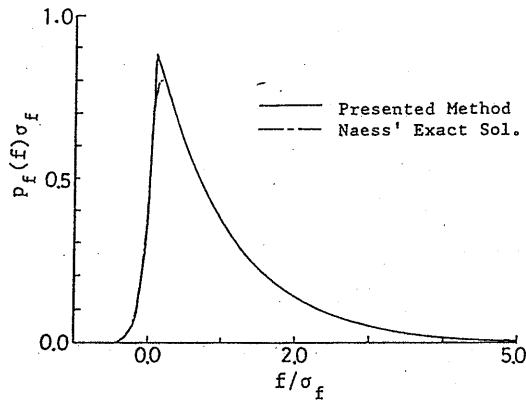
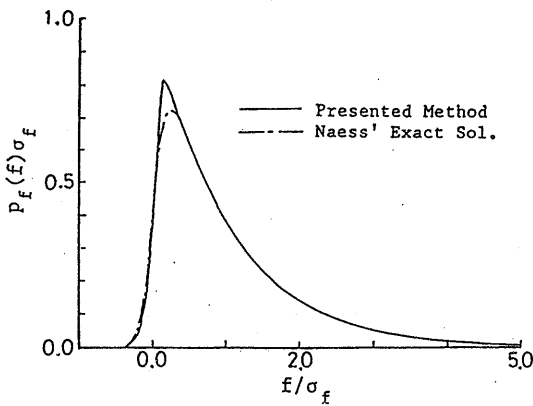


Figure 4.1 Explanatory sketch of a random process $X(t)$

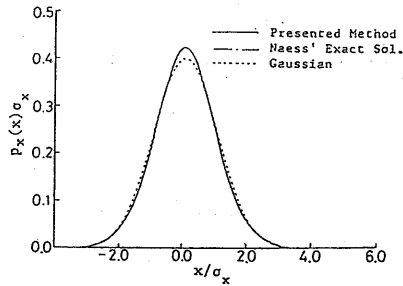


(a) Circular Structure

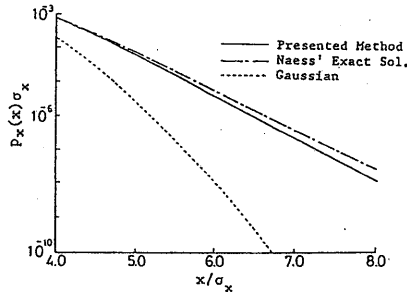


(b) Rectangular Structure

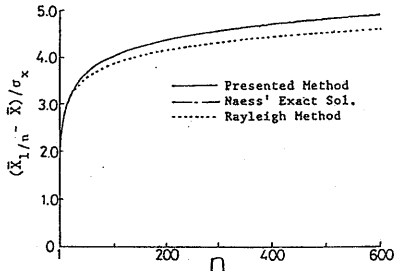
Figure 4.2 Instantaneous p.d.f. of pure second order forces



(a) P.D.F. of the Slowly Varying Sway Motion

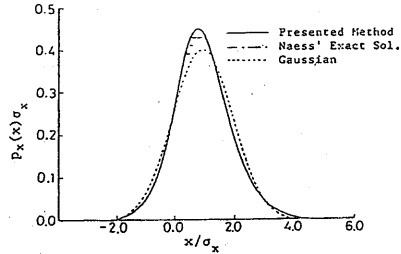


(b) P.D.F. Tail Behaviour

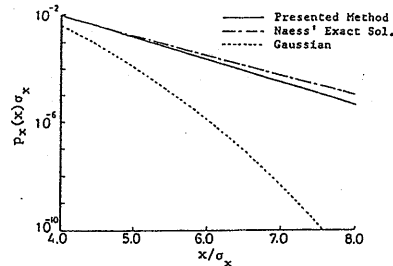


(c) 1/n th Highest Expected Amplitude

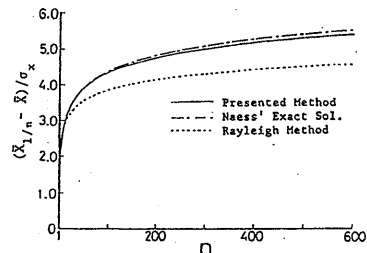
Figure 4.3 Instantaneous p.d.f. and 1/n th highest mean amplitude of pure second order responses (Case 1 - circular cylinder)



(a) P.D.F. of the Slowly Varying Sway Motion

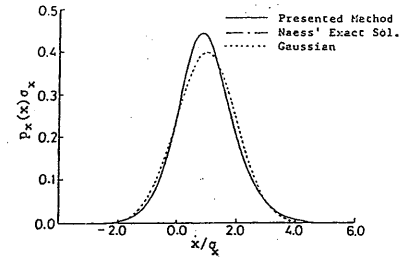


(b) P.D.F. Tail Behaviour

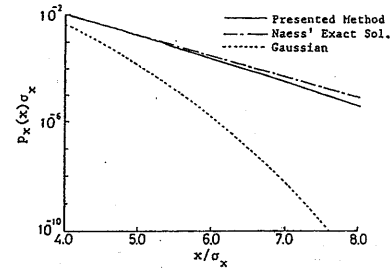


(c) 1/n th Highest Expected Amplitude

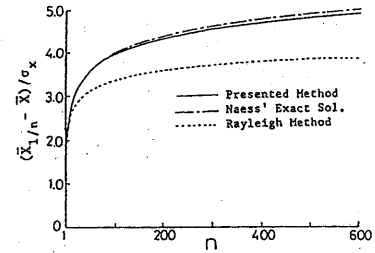
Figure 4.4 Instantaneous p.d.f. and 1/n th highest mean amplitude of pure second order responses (Case 2 - circular cylinder)



(a) P.D.F. of the Slowly Varying Sway Motion



(b) P.D.F. Tail Behaviour



(c) 1/n th Highest Expected Amplitude

Figure 4.5 Instantaneous p.d.f. and 1/n th highest mean amplitude of pure second order responses (Case 3 - rectangular cylinder)

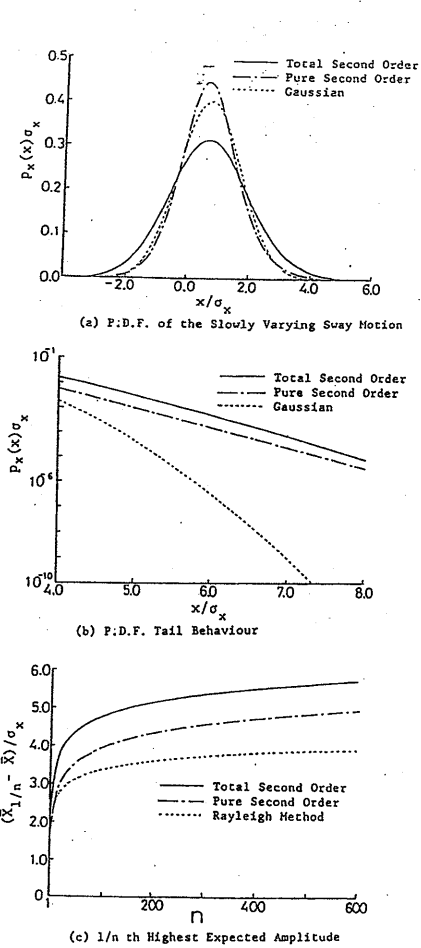


Figure 4.6 Statistical interference between first and second order responses for heavy damping ($\frac{\sigma_1}{\sigma_2} = 1.36$ and $\kappa = 0.1$)

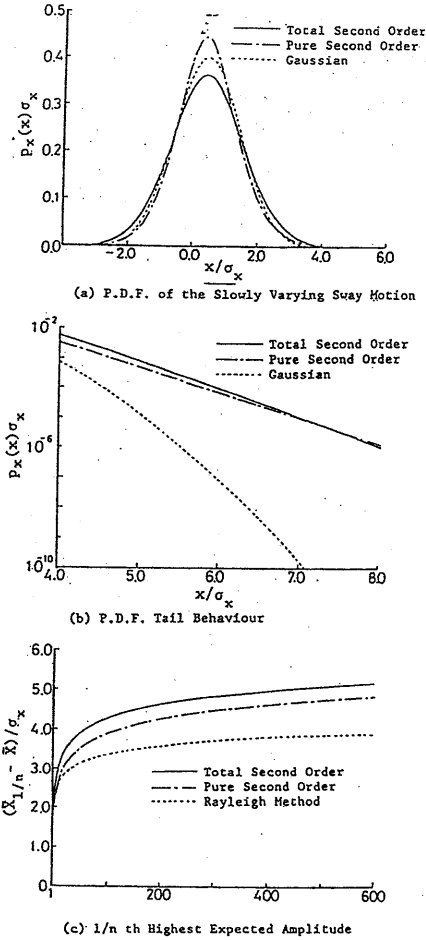


Figure 4.7 Statistical interference between first and second order responses for medium damping ($\frac{\sigma_1}{\sigma_2} = 2.9$ and $\kappa = 0.006$)

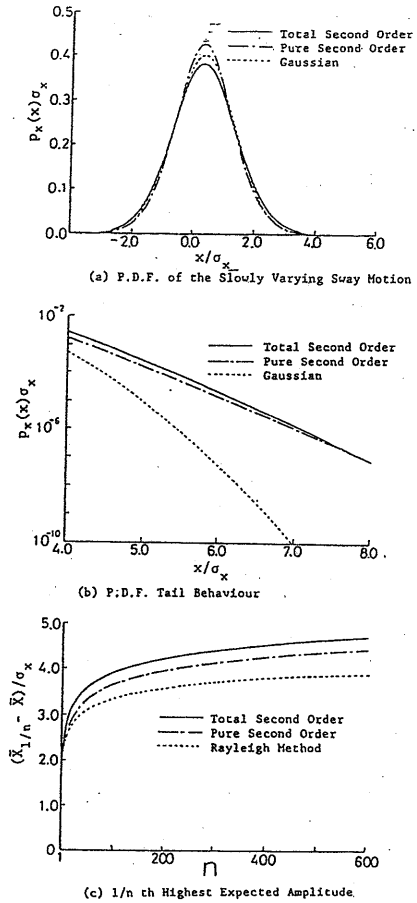
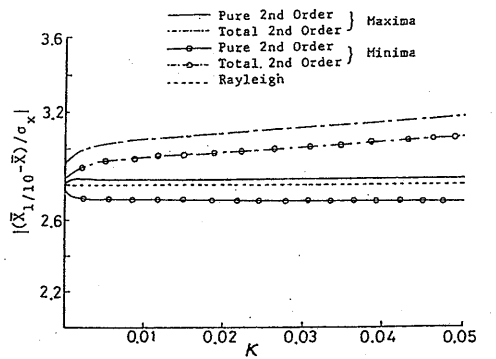
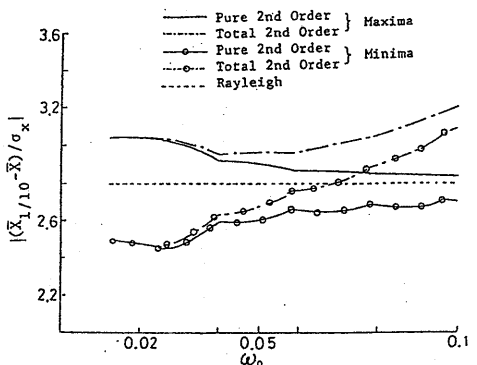


Figure 4.8 Statistical interference between first and second order responses for light damping ($\frac{\sigma_1}{\sigma_2} = 4.96$ and $\kappa = 0.0001$)

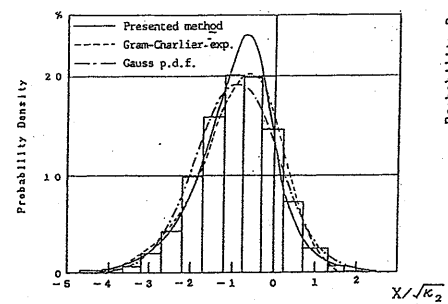


(a) 1/10th Highest Expected Amplitude VS. Damping Coefficient ($\omega_0=0.1$ rad/sec.)

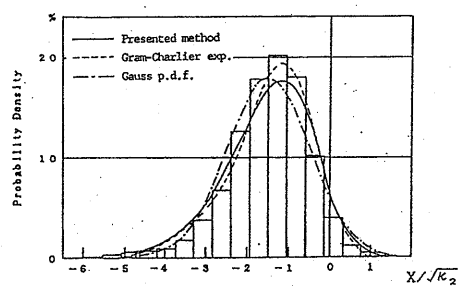


(b) 1/10th Highest Expected Amplitude VS. Natural Frequency ($\kappa=0.06$)

Figure 4.9 1/10 th highest mean amplitudes VS. damping coefficient and natural frequency

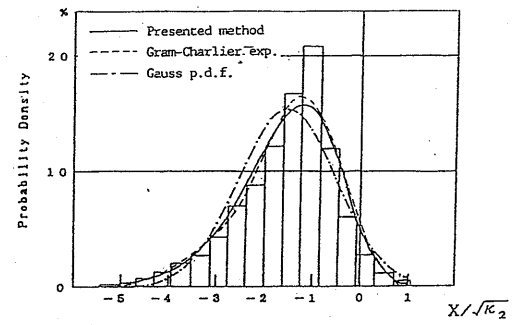


1) wave condition 1

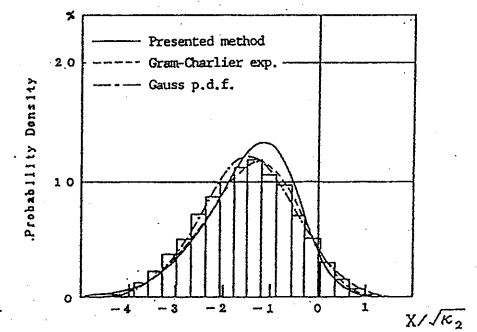


2) wave condition 2

Figure 4.10 Comparisons between observed histograms and estimated instantaneous probability density functions of surge motion (Wave conditions 1 and 2)



3) wave condition 3



4) wave condition 4

Figure 4.11 Comparisons between observed histograms and estimated instantaneous probability density functions of surge motion (Wave conditions 3 and 4)

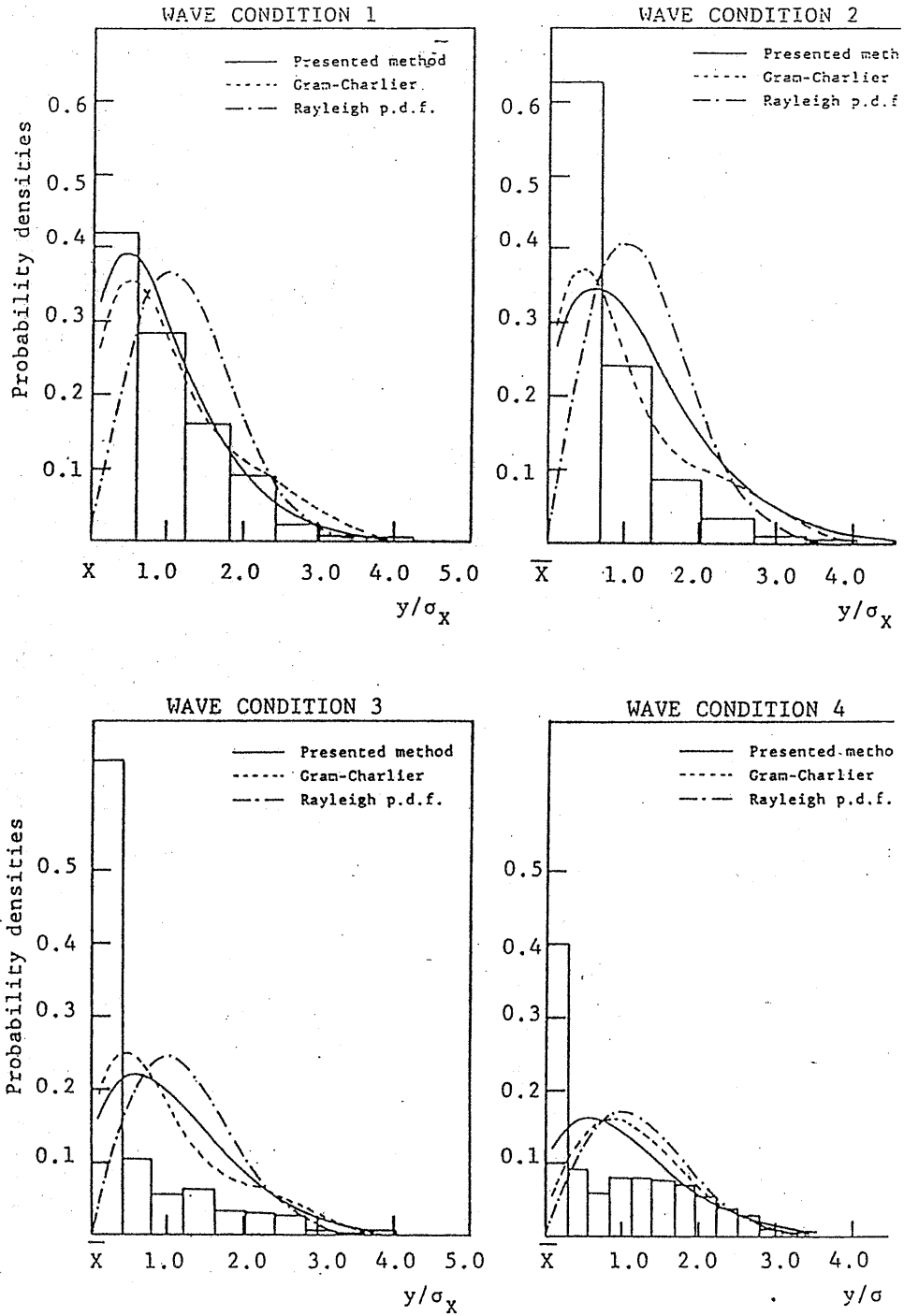


Figure 4.12 Comparisons between observed histograms and estimated maximum probability density functions of surge motion

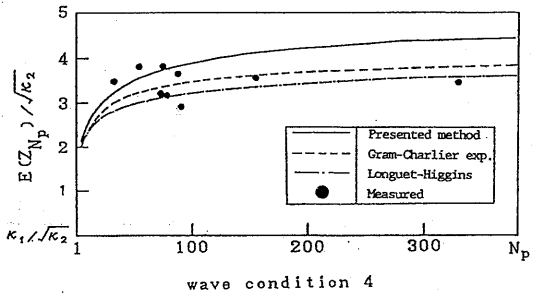
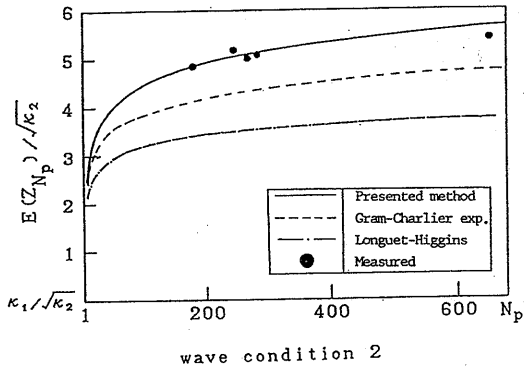
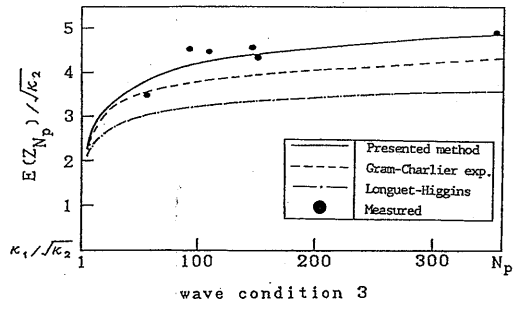
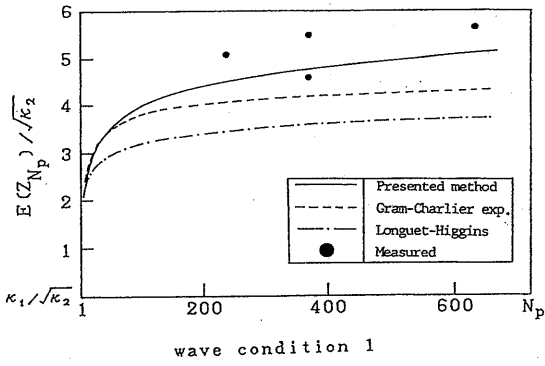


Figure 4.13 Comparisons between observed extreme responses and estimated ones (Wave conditions 1 and 2)

Figure 4.14 Comparisons between observed extreme responses and estimated ones (Wave conditions 3 and 4)

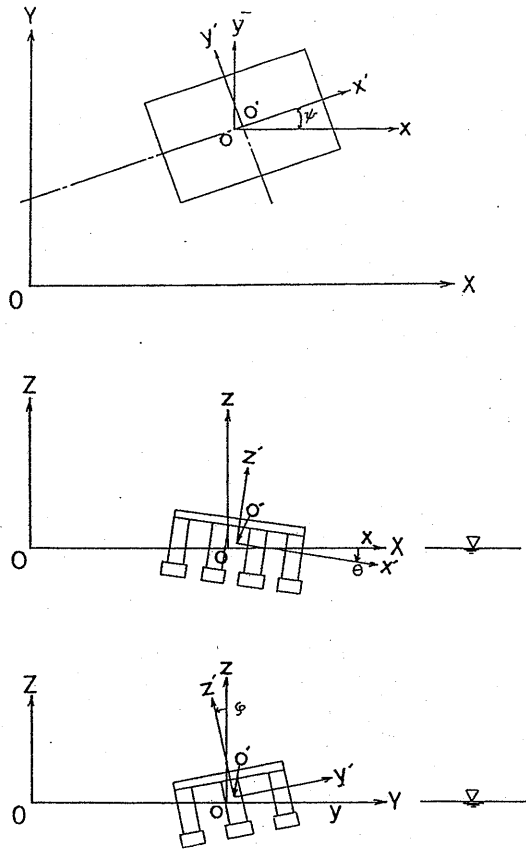


Figure A.1 System of coordinates

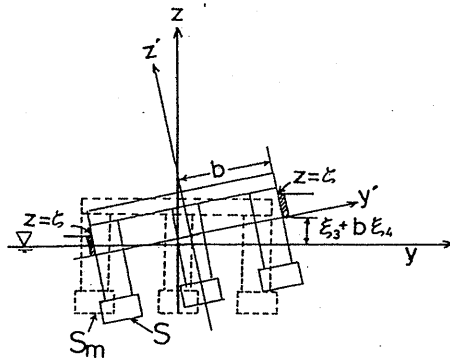


Figure A.2 Relationship between S and S_m

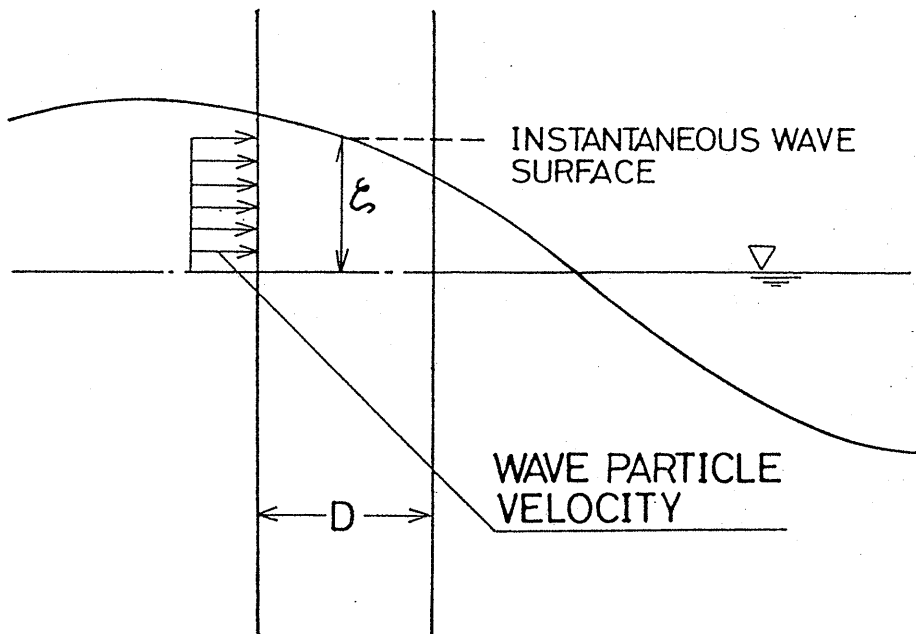


Figure C.1 Contribution to mean force from portion of wave between mean and instantaneous free surface